# SUSY effects in Kaon physics: Lepton Universality tests and rare decays

P. Paradisi

TU of Munich

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# General Considerations

## Flavor Physics in the LHC era

- High energy experiments are the key tool to determine the energy scale ∧ by direct production of NP particles.
- Low energy experiments are a fundamental ingredient to determine the symmetry properties of the new d.o.f. via their virtual effects in precision observables.

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General Considerations NP search strategies  $\mu - e$  universali

# General Considerations

G. Isidori - Flavour Physics now and in the LHC era

Flavour physics in the LHC era



LHC [high p<sub>T</sub>] A unique effort toward the

#### high-energy frontier



[to determine the energy scale of NP]

# Flavour physics



A *collective* effort toward the high-intensity frontier [to determine the <u>flavour structure</u> of NP]



# NP search strategies

Where to look for New Physics?

- Processes very suppressed or even forbidden in the SM
  - FCNC processes  $(\mu \to e\gamma, \tau \to \mu\gamma, B^0_{s,d} \to \mu^+\mu^-, K \to \pi\nu\bar{\nu})$
  - CPV effects (electron/neutron EDMs, *d<sub>e,n</sub>*....)
  - CPV in B<sub>s,d</sub> decay/mixing amplitudes
- Processes predicted with high precision in the SM
  - EWPO as  $\Delta \rho$ ,  $(g-2)_{\mu}$ ....
  - LU tests in  $R_M^{e/\mu} = \Gamma(K(\pi) \to e\nu) / \Gamma(K(\pi) \to \mu\nu)$

# LU tests with $K\ell$ 3 and $k\ell$ 2 modes

#### Conclusion and outlook

- The charged current analyses using K<sub>13</sub> and K<sub>12</sub> data have entered an era of very high precision
  - Improvements on the theoretical side: EM, isospin breaking corrections, dedicated dispersive parametrizations to analyse the FFs with the best precision.

On the experimental side, very precise data on K<sub>13</sub> and K<sub>12</sub> decays

Flavianet Kaon WG

- This allows for very precise tests of the SM (test of unitarity of the 1<sup>rst</sup> line of CKM matrix, universality, quark mass ratios...) and New Physics scenarios (Charged right-handed currents, scalar couplings, Lepton flavour violation...)
- But still on the experimental side, need K+ measurements (FFs..). Experimental puzzle on f<sub>0</sub>(t) (NA48 doesn't agree with the other experiments).
- On theoretical side, f<sub>+</sub>(0) determination should be improved
   disagreement between analytical and lattice determinations. Lattice improvements are promising.

E. Passemar

Kaon'09, Tsukuba

# $\mu - e$ universality in $R_K = \Gamma(K \to e \nu_e) / \Gamma(K \to \mu \nu_\mu)$

	$(R_{K}^{e/\mu})_{exp.} \ [10^{-5}]$
PDG 2006	$2.45\pm0.11$
NA48/2 '03 prel.	$2.416 \pm 0.043 \pm 0.024$
NA48/2 '04 prel.	$2.455 \pm 0.045 \pm 0.041$
KLOE prel.	$2.55 \pm 0.05 \pm 0.05$
SM prediction	$2.477\pm0.001$

$$({\it R}_{\it K}^{e/\mu})_{exp.}=(2.457\pm0.032) imes10^{-5}$$

• A dedicated run (of 4 month) for  $R_K$  by P326/NA62 (former NA48) has been performed @ the CERN. Goal: the error @ 0.3%!

• 
$$R_{\pi}^{exp.} = (1.230 \pm 0.004) \cdot 10^{-4}$$
 PDG

SUSY  $R_{K}^{LFV}$  in SUSY  $R_{K}^{LFV}$  in SUSY LU at a (Sup

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# SM prediction for $R_{K,\pi}$

• 
$$R_K^{SM} = (2.472 \pm 0.001) \cdot 10^{-5}$$
 SM

• 
$${\cal R}^{SM}_{\pi} = (1.2354 \pm 0.0002) \cdot 10^{-4}$$
 SM

Marciano Sirlin '93, Finkemeyer '96

• 
$$R_K^{SM} = (2.477 \pm 0.001) \cdot 10^{-5}$$
 SM

• 
$${\cal R}^{SM}_{\pi} = (1.2352 \pm 0.0001) \cdot 10^{-4}$$
 SM

Cirigliano Rossell '07

## The total errors in $R_{K,\pi}$ are dominated by the EXP. ERRORS!!!

# *R*<sub>*k*</sub> @ NA62



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<sup>V</sup> in SUSY  $R_K^{LFV}$  in SUSY  $R_K^{LFV}$  in SUSY LU at a (Sup

# R<sub>K</sub> @ KLOE



General Considerations NP search strategies  $\mu - e$  universal  $R_{K}^{LFV}$  in SUSY  $R_{K}^{LFV}$  in SUSY LU at a (Sup

# $\mu - e$ universality in $M \rightarrow I \nu$

• Any deviation from the SM expectation for  $R_{K,\pi}$  due to NP can be written as

$$R_{K,\pi} = R_{K,\pi}^{SM} \left( 1 + \Delta r_{K,\pi NP}^{e-\mu} \right),$$

- Violations of LU in CCI can be classified as
  - i) Corrections to  $(V-A) \times (V-A)$  interaction through  $W \ell \nu_{\ell}$  vertex correction induced by a loop of NP particles

$$\Delta r_{SUSY}^{e-\mu} \sim \frac{\alpha_2}{4\pi} \left( \frac{\tilde{m}_{\mu}^2 - \tilde{m}_{e}^2}{\tilde{m}_{\mu}^2 + \tilde{m}_{e}^2} \right) \frac{m_W^2}{M_{SUSY}^2} \leq 10^{-4}$$

• ii) New Lorentz Structures, i.e. scalar CCI with

 $H\ell\nu\sim m_\ell \tan\beta$ 

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# $\mu - e$ universality in $M \rightarrow l\nu$

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• Four-Fermi interaction for  $M \rightarrow \ell \nu$  induced by  $W^{\pm}$ ,  $H^{\pm}$ 

$$\frac{4G_F}{\sqrt{2}}V_{ud}\bigg[(\overline{u}\gamma_{\mu}P_Ld)(\overline{\ell}\gamma^{\mu}P_L\nu_{\ell})-t_{\beta}^2\bigg(\frac{m_dm_{\ell}}{m_{H^{\pm}}^2}\bigg)(\overline{u}P_Rd)(\overline{\ell}P_L\nu_{\ell})\bigg]$$

- PCAC's •  $< 0|\overline{u}\gamma_{\mu}\gamma_{5}d|M >= if_{M}p_{M}^{\mu}$   $< 0|\overline{u}\gamma_{5}d|M >= -if_{M}\frac{m_{M}^{2}}{m_{d}+m_{u}}$
- *H*<sup>±</sup> (**W**<sup>±</sup>) amplitude is proportional to *m*<sub>ℓ</sub> because of the Yukawa coupling (helicity suppression)

$$\frac{\Gamma^{H^{\pm}+W^{\pm}}(M\to\ell\nu)}{\Gamma^{W^{\pm}}(M\to\ell\nu)} = r_{\mathcal{M}} = \left[1 - t_{\beta}^{2} \left(\frac{m_{d}}{m_{u}+m_{d}}\right) \frac{m_{M}^{2}}{m_{H^{\pm}}^{2}}\right]^{2}$$

Tree level  $H^{\pm}$  effects  $(r_M)$  are lepton flavour blind.

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# $\mu - e$ universality in $M \rightarrow I\nu$

WHAT ARE WE MISSING ?.....

$$R_{K}^{EXP.} = \frac{\Gamma(\mathbf{K} \rightarrow \mathbf{e}\nu_{\mathbf{e}}) + \Gamma(\mathbf{K} \rightarrow \mathbf{e}\nu_{\mu}) + \Gamma(\mathbf{K} \rightarrow \mathbf{e}\nu_{\tau})}{\Gamma(\mathbf{K} \rightarrow \mu\nu_{\mu}) + \Gamma(\mathbf{K} \rightarrow \mu\nu_{e}) + \Gamma(\mathbf{K} \rightarrow \mu\nu_{\tau})}$$

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# $\mu - e$ universality in $M \rightarrow I\nu$

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# EXPERIMENTALLY THE NEUTRINO FLAVOUR IS **UNDETERMINED !!**

# $\mu - e$ universality in $M \rightarrow l\nu$

WHAT ARE WE MISSING ?.....

$$R_{K}^{\underline{\mathsf{EXP.}}} = \frac{\Gamma(\mathsf{K} \to \mathbf{e}\nu_{\mathbf{e}}) + \Gamma(\mathsf{K} \to \mathbf{e}\nu_{\mu}) + \Gamma(\mathsf{K} \to \mathbf{e}\nu_{\tau})}{\Gamma(\mathsf{K} \to \mu\nu_{\mu}) + \Gamma(\mathsf{K} \to \mu\nu_{e}) + \Gamma(\mathsf{K} \to \mu\nu_{\tau})}$$

# .....EXPERIMENTALLY THE NEUTRINO FLAVOUR IS UNDETERMINED !!

#### Masiero, Paradisi, Petronzio, '06

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# $R_{K}^{LFV}$ in SUSY



Which is the sign of  $\Delta r_{NP}^{e-\mu}$ ?

• LFV effects to LFC channels in  $R_M$ 

$${}^{\ell}H^{\pm}\nu_{\ell} \to \frac{g_2}{\sqrt{2}}\frac{m_{\ell}}{M_W}\tan\beta\left(1+\frac{m_{\tau}}{m_{\ell}}\Delta_{RL}^{\ell\ell}\tan\beta\right) \qquad (\ell=e,\mu)$$

$$\Delta_{RL}^{\ell\ell} \sim \frac{\alpha_1}{4\pi} \delta_{RR}^{\ell3} \delta_{LL}^{3\ell} f_{loop} \leq 10^{-4}$$

• Deviations from  $\mu - e$  universality in  $K_{l2}$  and  $\pi_{l2}$ 

$$\frac{R_{K,\pi}^{LFV}}{R_{K,\pi}^{SM}} \simeq \left[ \left( 1 - \frac{m_{\tau}}{m_e} \frac{m_{K,\pi}^2}{M_{H^{\pm}}^2} \Delta_{RL}^{11} \tan^3 \beta \right)^2 + \frac{m_{\tau}^2}{m_e^2} \frac{m_{K,\pi}^4}{M_{H^{\pm}}^4} |\Delta_R^{31}|^2 \tan^6 \beta \right]$$

 $R_{\rm k}^{\rm LFV} \simeq R_{\rm k}^{\rm SM}(1-0.032)$ ,

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# $R_{\kappa}^{LFV}$ in SUSY



Black points explain the  $(g-2)_{\mu}$  anomaly

Masiero, P.P., Petronzio, '08

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#### RLFV in SUSY



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 $R_{K}^{LFV}$  in SUSY  $R_{K}^{LFV}$  in SUSY  $R_{K}^{LFV}$  in SUSY LU at a (Sup

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# LU at a (Super)B factories

• 
$$R_{\tau}^{\mu/e} = \Gamma(\tau \to \mu \nu \bar{\nu}) / \Gamma(\tau \to e \nu \bar{\nu})$$
  
 $R_{\tau}^{\mu/e} \simeq 1 - 10^{-3} \left(\frac{t_{\beta}}{50}\right)^2 \left(\frac{200 \text{GeV}}{M_{H^{\pm}}}\right)^2$ 

#### Mursula et al. '83

• 
$$R_{B \to D}^{\tau/\ell} = \Gamma(B \to D\tau\nu) / \Gamma(B \to D\ell\nu)$$
  
Hou '92, Tanaka '95, Kiers & Soni '97

$$\frac{R_{B \to D}^{\tau/\mu}}{R_{B \to D}^{\tau/\mu}|_{SM}} \simeq 1 - 0.3 \left(\frac{t_{\beta}}{50}\right)^2 \left(\frac{200 \text{GeV}}{M_{H^{\pm}}}\right)^2$$

Nierste et al.'08, Kamenik & Mescia '08

# SUSY MFV scenario @ large tan $\beta$

#### How natural is the MFV SUSY scenario @ large $\tan \beta$ ?

- Top-Bottom Yukawa unification in GUT  $\Rightarrow \tan \beta = (m_t/m_b)$
- $m_h > 114 {
  m GeV}$  constraint better satisfied
- $\Delta a_{\mu} = (g-2)_{\mu}/2 = (3\pm1) imes 10^{-9}$  naturally explained
- WMAP constraints "naturally" satisfied Ellis et al.
- Correlations between BR(B  $\rightarrow \tau \nu$ ) and BR(B  $\rightarrow X_s \gamma$ ),  $\Delta M_{B_s}$ , BR(B<sub>s,d</sub>  $\rightarrow \ell^+ \ell^-$ ),  $(g - 2)_{\mu}$  and  $m_{h^0}$

#### Isidori, P.P., '06



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# Constraints

- $\mathbf{B} \rightarrow \mathbf{X}_{\mathbf{s}} \gamma$ :  $[1.01 < \mathbf{R}_{\mathbf{B}\mathbf{s}\gamma} < 1.24]$
- $\mathbf{a}_{\mu}$ :  $[2 < 10^{-9} (\mathbf{a}_{\mu}^{\exp} \mathbf{a}_{\mu}^{SM}) < 4]$
- $\mathbf{B} \to \mu^+ \mu^-$ :  $[\mathcal{B}^{exp} < 8.0 \times 10^{-8}]$
- $\Delta M_{B_s}$  :  $[\Delta M_{B_s} = 17.35 \pm 0.25 \text{ ps}^{-1}]$

•  $B \to \tau \nu$  :  $[0.8 < \mathbf{R}_{\mathbf{B}\tau\nu} < 0.9]$ 

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# B-physics & $(g-2)_{\mu}$ under WMAP constraints



# $K \rightarrow \pi \nu \bar{\nu}$ in the SM

- $K \to \pi \nu \overline{\nu}$  processes offer a unique possibility in probing the underlying flavour mixing mechanism:
  - No SM tree-level contributions (FCNC decays );
  - One-loop SM contributions CKM-suppressed ( $V_{ts}^* V_{td} \sim \lambda^5$ );
  - High precision of the SM prediction thanks to short distance (e.w.) dynamics dominance:

$$\mathcal{H}_{\text{eff}}^{(s.d.)} = \sum_{l=e,\mu,\tau} V_{ts}^* V_{td} \left[ X_L(\bar{s}d)_{V-A} + X_R(\bar{s}d)_{V+A} \right] (\bar{\nu}_l \nu_l)_{V-A}$$

$$Br(K \to \pi \nu \bar{\nu}) \sim (X = X_L + X_R)^2$$

$$X \sim c_{SM} \frac{y_t^2 V_{ts}^* V_{td}}{16\pi^2 M_W^2}$$

$$X_L^{\text{SM}} = 1.464 \pm 0.041 , \quad X_R^{\text{SM}} = 0$$

$$V_L^{\text{SM}} = 0$$

# $K \rightarrow \pi \nu \bar{\nu}$ and NP

K → πνν has a high sensitivity to NP effects of many theories as SUSY, LHT, Z' models....

$$X(s 
ightarrow d)_{
m FCNC} \sim c_{
m SM} rac{y_t^2 V_{ts}^* V_{td}}{16 \pi^2 M_W^2} + c_{
m NP} rac{\delta_{
m 21}}{16 \pi^2 \Lambda_{NP}^2}$$

$$Br(K \to \pi \nu \overline{\nu}) \sim (X = X_L + X_R)^2$$

$$X_L = X_L^{\text{SM}} + X_L^{\text{NP}}, \qquad \qquad X_R = X_R^{\text{NP}}$$

• Large **NP** effects only if  $\delta_{21} \approx V_{ts}^* V_{td}$  (beyond MFV)

see the talk by Buras

K G. Isidori – Exploring BSM with K physics

Flavour in the era of the LHC

#### <u>Rare Kaon decays beyond the SM</u> [general properties]

Two basic scenarios:

Minimal Flavour Violation

flavour symmetry broken only by the (SM) Yukawa couplings

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Small deviations (10-20%) from SM

•Stringent correlations with other rare decays in B physics  $[B_d \rightarrow X_{s,d}vv,$ 

 $B_d \rightarrow X_{s,d} l^+ l^-, B_{s,d} \rightarrow l^+ l^-$ ]

Recent (almost) model-indep. analysis :



Consistent with results of specific models:

- Constrained MSSM [Buras et al. '01]
- One universal extra dim. [Buras et al. '03]
- Littlest-Higgs [Buras et al. '05]

# Gaugino mediated $K \rightarrow \pi \nu \bar{\nu}$

- The dominant effects to  $K \to \pi \nu \bar{\nu}$  arise from  $\tilde{\chi}/\tilde{u}$  diagrams with double-MIA [Colangelo, Isidori '98].
- Gluino-type amplitudes (LL, RR and LR-down squarks type mixings) essentially negligible contrary to  $\epsilon_K$ ,  $b \rightarrow s\gamma$ ,  $B^0 - B^0$
- Minor effects within pure MFV.
- The maximal sensitivity to the up-type trilinear terms is obtained for
  - Light stop and charginos
  - small tan  $\beta$



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General Considerations NP search strategies  $\mu - e$  universal

# $K \rightarrow \pi \nu \bar{\nu}$ in SUSY beyond MFV

G. Isidori - Exploring BSM with K physics

Flavour in the era of the LHC

#### <u>Rare Kaon decays beyond the SM</u> [general properties]

Two basic scenarios:

#### E.g.: II. Generic MSSM



New sources of Flavour Symmetry

breaking around the TeV scale

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- •Potentially large effects, especially in the three CPV  $K_L$  decays (no  $\lambda^5$  suppression)
- Correlations with observables in B physics not obvious

#### Grossman-Nir bound:

 $\frac{\mathcal{A}}{10} \qquad \Gamma(K_L \to \pi^0 \text{VV}) < \Gamma(K^+ \to \pi^+ \text{VV}) \qquad \geq \\ \text{SUSY effects in Kaon physics: Lepton Universality tests and rates and rates are supported as the superscript states are supported as the superscript state states are supported as the superscript state state state states are supported as the superscript states are supported as the superscript state state state states are supported as the superscript state state state states are superscript states are supported as the superscript state state state states are superscript states are super$ 

# Chargino mediated $K \rightarrow \pi \nu \bar{\nu}$



G.Isidori, F.Mescia, P.P., C.Smith, S.Trine, '06

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# Chargino mediated $K \rightarrow \pi \nu \bar{\nu}$



G.Isidori, F.Mescia, P.P., C.Smith, S.Trine, '06

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# Chargino mediated $K \rightarrow \pi \nu \bar{\nu}$



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# Higgs mediated $K \rightarrow \pi \nu \bar{\nu}$



# Conclusion

#### Where to look for New Physics?

- LU breaking @ % in  $R_K^{e/\mu}$  from SUSY LFV effects
- LU breaking @ % in  $R_{\kappa}^{e/\mu}$  implies  $M_{H^{\pm}} < 1 \text{TeV}$  and it can be compatible with the  $(g-2)_{\mu}$  anomaly
- LU breaking @ % in  $R_{\kappa}^{e/\mu} \Rightarrow BR(\tau \to e\gamma) > 10^{-9}$
- The relevant SUSY parameter space for  $R_K^{e/\mu}$  @ % is allowed by the constraints of rare LFV decays, **B**-physics observables and **D**ark Matter



#### Conclusion

# Conclusion

#### Where to look for New Physics?

- $K \rightarrow \pi \nu \overline{\nu}$  is a **golden channel** where to look for **NP** because:
  - $K \to \pi \nu \overline{\nu}$  is predicted with high resolution in the SM
  - $K \rightarrow \pi \nu \overline{\nu}$  has a **high sensitivity to NP** effects of many theories as SUSY, LHT, Z' models....

## ₩

## Kaon Physics will play a major role in the LHC ERA to unveil and to understand NP

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